

NNLO QCD predictions for single jet inclusive production at the LHC

J. Currie^a, E.W.N. Glover^a, J. Pires^b

^a *Institute for Particle Physics Phenomenology, University of Durham, Durham DH1 3LE, England*

^b *Max-Planck-Institut für Physik, Föhringer Ring 6 D-80805 Munich, Germany*

We report the first calculation of fully differential jet production in all partonic channels at next-to-next-to leading order (NNLO) in perturbative QCD and compare to the available ATLAS 7 TeV data. We discuss the size and shape of the perturbative corrections along with their associated scale variation across a wide range in jet transverse momentum, p_T , and rapidity, y . We find significant effects, especially at low p_T , and discuss the possible implications for Parton Distribution Function fits.

PACS numbers: 13.87.Ce, 12.38.Bx

The Large Hadron Collider (LHC) is currently colliding protons at centre of mass energies up to $\sqrt{s} = 13$ TeV. The main goal is to search the high energy frontier for signs of physics beyond the Standard Model. However, any searches for new physics are irreducibly dependent on how well we understand the Standard Model and the collider environment of the LHC itself.

At the LHC the inclusive cross section for a given final-state can be calculated using the factorization formula,

$$d\sigma = \sum_{i,j} \int \frac{d\xi_1}{\xi_1} \frac{d\xi_2}{\xi_2} f_i(\xi_1, \mu_F) f_j(\xi_2, \mu_F) d\hat{\sigma}_{ij} \quad (1)$$

which is accurate up to non-perturbative hadronization corrections, typically of the order Λ_{QCD}/Q , where Q is the hard scale in the scattering process. The partonic cross section, $d\hat{\sigma}_{ij}$, can be calculated as a perturbative series in the strong coupling, α_s , and systematically improved by progressively including higher order terms in the series. It is also necessary to have a good understanding of the non-perturbative Parton Distribution Functions (PDF), $f_i(\xi, \mu_F)$. The PDFs quantify the relative parton content of the proton carrying a fraction, ξ , of the proton's momentum for a given factorization scale, μ_F . To calculate the cross section using this formula we need accurate determinations of the PDFs, α_s and the higher order terms in the perturbative expansion of the partonic cross section.

Data from lepton-nucleon Deep Inelastic Scattering (DIS) experiments such as HERA [1] provide detailed information about the quark PDFs and have been used to significantly constrain the uncertainties on these quantities. The inclusive cross section in DIS involves the exchange of a virtual photon coupling to quarks at lowest order via the electroweak coupling constant. The electrical neutrality of the gluon means that the gluon PDF can only be constrained using specific final-states, such as heavy quarks or jets [2], or indirectly through DGLAP evolution of the flavour singlet distribution. In contrast, jet production at the Tevatron [3, 4] and LHC directly probes the gluon PDF and is $\mathcal{O}(\alpha_s^2)$ at leading order (LO). The single jet inclusive cross section has been

measured accurately by ATLAS [5, 6] and CMS [7] across the large dynamical range of the LHC.

To take advantage of the available data we must be able to calculate observables with sufficient precision yet the cross section for producing jets is currently only known exactly at next-to leading order (NLO) [8–12] and partially at NNLO [13]. The theoretical uncertainty in this observable, estimated from the dependence on unphysical scales, is the main limiting factor when determining parameters like α_s from jet data or consistently including this data in global fits for PDFs [14–17]. To improve on the status-quo it is clear that an accurate and precise determination of jet production at the LHC is needed and so in this letter we present the first calculation of the NNLO correction to jet production in perturbative QCD. Higher order corrections have the potential to change the size and shape of the cross section and also to reduce the residual scale dependence in a calculation; we discuss the extent to which this is true for the NNLO correction to the fully differential single jet inclusive cross section.

Predictions for jet production at NNLO accuracy require the relevant tree-level [18], one-loop [19–21] and two-loop [22–24] parton-level scattering amplitudes as well as a procedure for dealing with the infrared (IR) singularities present in both the phase space integrals and matrix elements, but which cancel in any IR safe physical observable. Several techniques have been developed for obtaining finite cross sections at NNLO for hadronic initial-states: antenna subtraction [25, 26], q_T -subtraction [27], N -jettiness subtraction [28], sector-improved residue subtraction [29], sector decomposition [30] and projection to Born [31]. We use the antenna subtraction method, implemented in the parton-level event generator, NNLOJET [32, 33], to calculate the single jet inclusive cross section, fully differential in the jet transverse momentum, p_T and rapidity, y .

We include the leading colour contribution from all partonic subprocesses in all channels. For example, in the gluon-gluon scattering channel there are three partonic subprocesses contributing to the double real correction: $gg \rightarrow gggg$, $gg \rightarrow q\bar{q}gg$ and $gg \rightarrow q\bar{q}q\bar{q}$; we include the contributions which are leading in the number of colours,

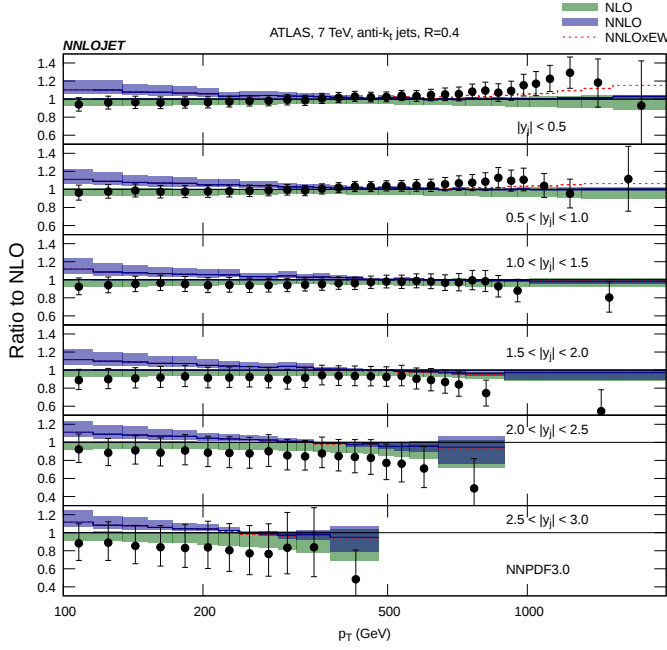


FIG. 1: Double-differential inclusive jet cross-sections measurement by ATLAS [6] and NNLO perturbative QCD predictions as a function of the jet p_T in slices of rapidity, for anti- k_T jets with $R = 0.4$ normalized to the NLO result. The shaded bands represent the scale uncertainty of the theory predictions obtained by varying μ_R and μ_F as described in the text. The red dashed line displays the NNLO/NLO ratio corrected multiplicatively for electroweak corrections [37].

N_c , to all these subprocesses. In practice this amounts to calculating the N_c^2 , $N_c N_F$ and N_F^2 corrections to all LO subprocesses, where N_F is the number of light quark flavours. We include the full LO and NLO coefficients in this calculation but note that retaining only the leading colour correction to all partonic subprocesses at NLO gives the full result to within a few percent across all distributions. The analogous subleading colour contributions at NNLO are expected to be small and we do not include them in this study. To support this assumption we note that the subleading colour NNLO contribution for pure gluon scattering was presented in a previous study [34] and found to be negligible. We construct subtraction terms to regulate all IR divergences in the phase space integrals and cancel all explicit poles in the dimensional regularization parameter, $\epsilon = (4 - d)/2$, the details of which for the antenna subtraction method can be found in [25, 34, 36]. The IR finite cross section at NNLO is then integrated numerically in four dimensions over the appropriate two-, three- or four-parton massless phase space to yield the final result.

In Fig. 1 we present the results for the double-differential inclusive jet cross section at NLO and NNLO, normalized to the NLO theoretical prediction to emphasize the impact of the NNLO correction to the NLO re-

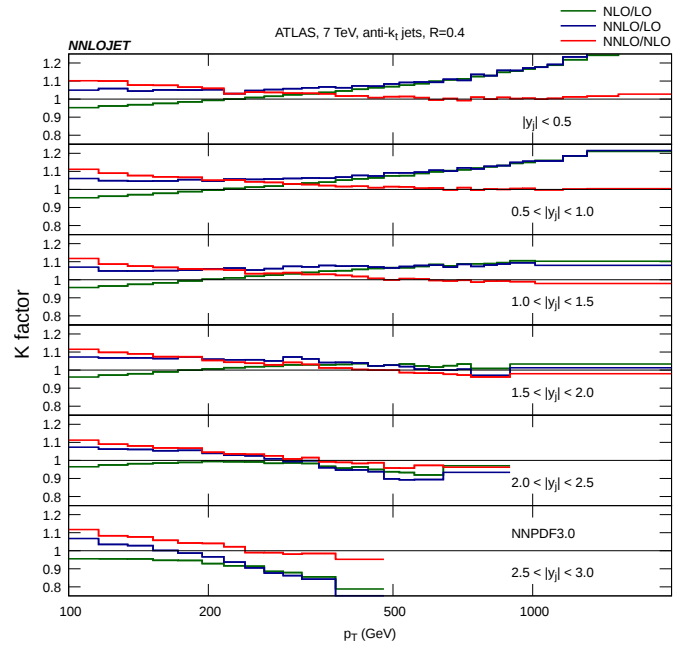


FIG. 2: NLO and NNLO k -factors for jet production at $\sqrt{s} = 7$ TeV. The lines correspond to the double differential k -factors (ratios of perturbative predictions in the perturbative expansion) for $p_T > 100$ GeV and across six rapidity $|y|$ slices.

sult. The collider setup is proton-proton collisions at a centre of mass energy of $\sqrt{s} = 7$ TeV where the jets are reconstructed using the anti- k_T jet algorithm [35] with $R = 0.4$. We use the NNPDF3.0 NNLO PDF set [15] with $\alpha_s(M_Z^2) = 0.118$ throughout this paper for LO, NLO and NNLO predictions to emphasise the behaviour of the higher order coefficient functions at each perturbative order. By default we set the renormalization and factorization scales $\mu_R = \mu_F = p_{T1}$, where p_{T1} is the p_T of the leading jet in each event. To obtain the scale uncertainty of the theory prediction we vary both scales independently by a factor of $1/2$ and 2 with the constraint $1/2 \leq \mu_R/\mu_F \leq 2$. We find that the NNLO coefficient has a moderate positive effect on the cross section, 10% at low p_T across all rapidity slices relative to NLO. This is significant because it is precisely in this region where the majority of the cross section lies, especially in the central rapidity slices, and it is where we observe the largest NNLO effects. At higher p_T we see that the relative size of the NNLO correction to NLO decreases to the 1-2% level and so the perturbative series converges rapidly.

Given that we see a moderate NNLO correction to the NLO prediction in the region where the bulk of the cross section lies, it is instructive to compare to the available data. The data points in Fig. 1 represent the ATLAS data for an integrated luminosity of 4.5 fb^{-1} [6], normalized to the NLO prediction. We do not include non-

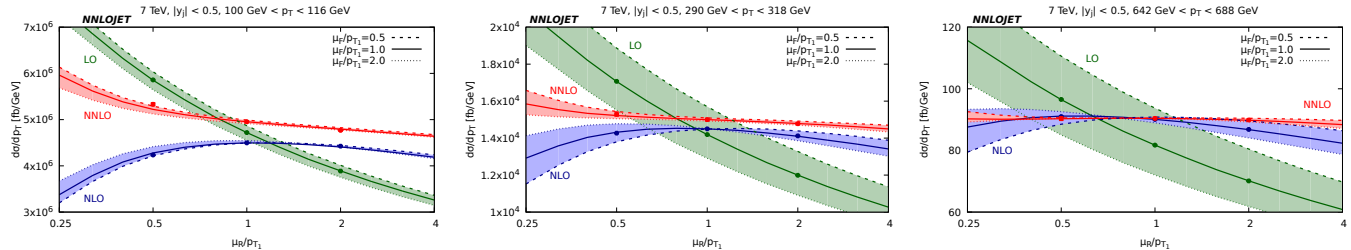


FIG. 3: The single inclusive cross section in three p_T bins: (left) 100-116 GeV, (centre) 290-318 GeV, (right) 642-688 GeV, plotted against the normalized scale choice μ_R/p_{T1} . Points represent the LO, NLO and NNLO cross section as computed by NNLOJET at $\mu_R/p_{T1} = 0.5, 1, 2$ and $\mu_F/p_{T1} = 1$. The solid lines represent the Renormalization Group Equation (RGE) solution for the scale variation, computed with p_T as the evolution parameter and $\mu_F/p_{T1} = 1$. Long and short dashed lines represent the same quantities evaluated with $\mu_F/p_{T1} = 0.5$ and $\mu_F/p_{T1} = 2$ respectively.

perturbative effects in our predictions; they are quantified in [6] and found to be a 2% effect in the lowest p_T bin and at most a 1% effect in all other bins (although the quoted uncertainty on the non-perturbative corrections can be as high as 9% for the lowest p_T bin). The electroweak corrections computed in [37], are applied multiplicatively to the QCD calculation for central scale choice using the information provided in [6] and the total is displayed as the red dashed line. The electroweak effects are small to moderate for $p_T > 1$ TeV for central rapidities but otherwise negligible. We observe that the data is described very accurately by the NLO prediction, particularly at low to moderate p_T , whilst the NNLO prediction shows some tension with the data in the same region.

The potential for the NNLO correction to change the shape of the distribution relative to NLO can be seen clearly in Fig. 2 where we show the k -factors for NLO/LO, NNLO/NLO and NNLO/LO as a function of p_T in six rapidity slices. In the central rapidity slices we observe that at low p_T the NLO correction acts negatively relative to LO and then grows to a moderate positive correction at high p_T . In contrast, the NNLO correction acts positively at low p_T and decreases to a small effect at high p_T . The aggregate effect is shown in the NNLO/LO curve which is the result of a partial cancellation between NLO and NNLO at low p_T , reversing the negative NLO contribution to give a positive total correction, and largely follows the NLO curve at high p_T . In the region, $|y| > 1.5$, we observe that the NLO correction is once again negative at low p_T but does not grow as strongly as in central regions, and is indeed negative at high p_T for the most forward slices. Relative to NLO, the NNLO corrections are again positive and moderate at low p_T , decreasing in size at high p_T , such that the total effect is a positive correction at low p_T decreasing to a negative effect at high p_T in the most forward region.

Aside from the size and shape of the NNLO corrections, an interesting feature of Fig. 1 is the scale dependence at NLO and NNLO, represented by the thickness of the

bands. At high p_T , especially in the central region, the NNLO scale dependence is dramatically reduced and the NNLO band lies firmly within the NLO band. The situation is once again different at low p_T where we observe an appreciable scale variation on the NNLO calculation, in some places even larger than the NLO scale variation, and the bands do not fully overlap.

This behaviour is unexpected and so in Fig. 3 we analyse the scale variation in more detail. We select three p_T bins at low (100-116 GeV), intermediate (290-318 GeV) and high (642-688 GeV) p_T in the central region $|y| < 0.5$ and display the cross section as a function of μ_R/p_{T1} . The points represent the cross section as calculated at LO, NLO and NNLO, evaluated at $\mu_R/p_{T1} \in \{0.5, 1, 2\}$ with $\mu_F/p_{T1} = 1$. The solid line is the Renormalization Group Equation (RGE) prediction for the scale variation starting from the cross section computed with $\mu_R/p_{T1} = \mu_F/p_{T1} = 1$. To be fully consistent, the RGE evolution variable should be p_{T1} , however starting with the distribution $d\sigma/dp_T$, all information on p_{T1} is lost and therefore we use p_T to approximate p_{T1} as the RGE evolution variable [38]. Nevertheless, the difference in the evolution is small even at low p_T and we include the RGE lines to aid the discussion of the scale variation in each bin. The long and short dashed lines are obtained using $\mu_F/p_{T1} = 0.5$ and $\mu_F/p_{T1} = 2$ respectively.

From the left pane in Fig. 3 we observe that at LO the scale variation is a monotonically decreasing function with p_T . At NLO the picture is quite different; the shape of the RGE curve turns over at approximately the central scale choice. The consequence of this behaviour is that the scale band is asymmetric, with the central scale being located at the upper edge of the band, as can be seen in Fig. 1 where the NLO scale band lies almost entirely below one. The overall variation is also relatively small at low p_T , ($< 5\%$), which is linked to the smallness of the NLO coefficient, ($\sim 4\%$ of LO), as displayed in Fig. 2. The NNLO curve is monotonically decreasing, leading to a more symmetric band and the overall variation is significant, largely reflecting the size of the

NNLO correction. In the central and right panes of Fig. 3 we observe a similar behaviour; the NLO curve is once again turning over, yielding an asymmetric scale band in Fig. 1. The NNLO curves are monotonically decreasing and show diminishing scale variation with increasing p_T as the NNLO coefficient becomes negligible relative to the moderately large NLO correction, as displayed in Fig. 2.

To understand these results it is instructive to consider the low and high p_T regions separately. At high p_T , the NLO correction is moderate and positive, the NNLO correction is small and reduces the scale variation significantly relative to NLO with the NNLO scale band lying within the NLO band. At low p_T the behaviour is quite different; we observe that the NLO correction is small and generally negative whereas the NNLO correction is moderate and positive. The relative size of the NLO and NNLO corrections leads to a significant NNLO scale band which largely does not overlap with the NLO band.

A possible explanation may be found in the comparison to data in Fig. 1. The data fits the NLO prediction very closely when using the NNLO NNPDF3.0 PDF set. This PDF set is fitted to Tevatron and LHC jet data at NNLO despite the exact calculation for jet production being reported for the first time in this letter. In the approximate NNLO PDF fit, a restricted data set at high p_T was used in conjunction with an approximation to the NNLO coefficient functions. This procedure has a small influence on the low p_T region which was included in the NLO PDF fit and in fact we find that the NLO prediction using the NLO PDF set gives an almost identical result to that shown in Fig. 1 obtained with NNLO PDFs. However, at low p_T we see that the exact NNLO correction is not small and the NNLO PDF fit potentially underestimates the effect of the NNLO contribution in that region.

We have presented the first calculation of the single jet inclusive cross section including all partonic subprocesses at NNLO for the LHC using the antenna subtraction method, implemented in the parton-level event generator NNLOJET and compared to ATLAS data. We find that the NNLO corrections are moderate at low p_T and change the shape of the distributions relative to NLO. At high p_T the corrections are smaller and we see a dramatic reduction in the scale variation. The fact that the NNLO corrections move the theoretical prediction away from the available data suggests that this calculation may have an appreciable impact when included in refitting the PDF used in this study to jet data. We anticipate that this calculation will open the door for precision phenomenology using LHC jet data, including studies of scale choice, jet shape, cone size and different PDF sets.

The authors thank Xuan Chen, Juan Cruz-Martinez, Aude Gehrmann-De Ridder, Thomas Gehrmann, Alexander Huss, Tom Morgan and Jan Niehues for useful discussions and their many contributions to the

NNLOJET code. We thank Johannes Bluemlein, Joey Huston and Klaus Rabbertz for their useful comments on the draft manuscript. We gratefully acknowledge the assistance provided by Jeppe Andersen utilizing the computing resources provided by the WLCG through the GridPP Collaboration. This research was supported in part by the UK Science and Technology Facilities Council, in part by the Research Executive Agency (REA) of the European Union under the Grant Agreement PITN-GA-2012-316704 (“HiggsTools”) and the ERC Advanced Grant MC@NNLO (340983).

-
- [1] R. Devenish and A. Cooper-Sarkar, *Deep inelastic scattering*, Oxford University Press (Oxford, 2004).
 - [2] P. Newman and M. Wing, *Rev. Mod. Phys.* **86** (2014) 1037.
 - [3] A. Abulencia *et al.* [CDF Collaboration], *Phys. Rev. D* **75**, 092006 (2007) Erratum: [*Phys. Rev. D* **75**, 119901 (2007)].
 - [4] V. M. Abazov *et al.* [D0 Collaboration], *Phys. Rev. D* **85**, 052006 (2012)
 - [5] G. Aad *et al.* [ATLAS Collaboration], *Eur. Phys. J. C* **71**, 1512 (2011); *Eur. Phys. J. C* **73**, 2509 (2013).
 - [6] G. Aad *et al.* [ATLAS Collaboration], *JHEP* **1502**, 153 (2015) Erratum: [*JHEP* **1509**, 141 (2015)].
 - [7] V. Khachatryan *et al.* [CMS Collaboration], *Eur. Phys. J. C* **76**, 451 (2016); *JHEP* **1206**, 036 (2012); *Phys. Rev. Lett.* **107**, 132001 (2011).
 - [8] S. D. Ellis, Z. Kunszt and D. E. Soper, *Phys. Rev. Lett.* **69** (1992) 1496.
 - [9] W. T. Giele, E. W. N. Glover and D. A. Kosower, *Phys. Rev. Lett.* **73** (1994) 2019.
 - [10] Z. Nagy, *Phys. Rev. Lett.* **88** (2002) 122003; *Phys. Rev. D* **68** (2003) 094002.
 - [11] S. Alioli, K. Hamilton, P. Nason, C. Oleari and E. Re, *JHEP* **1104** (2011) 081.
 - [12] J. Gao, Z. Liang, D. E. Soper, H. L. Lai, P. M. Nadolsky and C.-P. Yuan, *Comput. Phys. Commun.* **184**, 1626 (2013).
 - [13] A. Gehrmann-De Ridder, T. Gehrmann, E. W. N. Glover and J. Pires, *Phys. Rev. Lett.* **110**, no. 16, 162003 (2013); J. Currie, A. Gehrmann-De Ridder, T. Gehrmann, N. Glover, J. Pires and S. Wells, arXiv:1407.5558 [hep-ph].
 - [14] S. Alekhin, J. Blumlein and S. Moch, *Phys. Rev. D* **89**, 054028 (2014)
 - [15] R. D. Ball *et al.* [NNPDF Collaboration], *JHEP* **1504**, 040 (2015).
 - [16] L. A. Harland-Lang, A. D. Martin, P. Motylinski and R. S. Thorne, *Eur. Phys. J. C* **75**, 204 (2015).
 - [17] S. Dulat *et al.*, *Phys. Rev. D* **93**, 033006 (2016).
 - [18] M. L. Mangano and S. J. Parke, *Phys. Rept.* **200** (1991) 301.
 - [19] Z. Bern, L. J. Dixon and D. A. Kosower, *Phys. Rev. Lett.* **70**, 2677 (1993).
 - [20] Z. Bern, L. J. Dixon and D. A. Kosower, *Nucl. Phys. B* **437**, 259 (1995).
 - [21] Z. Kunszt, A. Signer and Z. Trocsanyi, *Phys. Lett. B* **336**, 529 (1994).

- [22] E. W. N. Glover, C. Oleari and M. E. Tejeda-Yeomans, Nucl. Phys. B **605** (2001) 467; E. W. N. Glover and M. Tejeda-Yeomans, JHEP **0105** (2001) 010; Z. Bern, A. De Freitas and L. J. Dixon, JHEP **0203** (2002) 018.
- [23] C. Anastasiou, E. W. N. Glover, C. Oleari and M. E. Tejeda-Yeomans, Nucl. Phys. B **605**, 486 (2001); Phys. Lett. B **506**, 59 (2001); JHEP **0306**, 033 (2003); Z. Bern, A. De Freitas and L. J. Dixon, JHEP **0306**, 028 (2003). Erratum: [JHEP **1404**, 112 (2014)].
- [24] C. Anastasiou, E. W. N. Glover, C. Oleari and M. E. Tejeda-Yeomans, Nucl. Phys. B **601**, 318 (2001); Nucl. Phys. B **601**, 341 (2001); Phys. Lett. B **506**, 59 (2001); A. De Freitas and Z. Bern, JHEP **0409**, 039 (2004).
- [25] A. Gehrmann-De Ridder, T. Gehrmann and E. W. N. Glover, JHEP **0509**, 056 (2005); J. Currie, E. W. N. Glover and S. Wells, JHEP **1304**, 066 (2013).
- [26] A. Daleo, T. Gehrmann and D. Maitre, JHEP **0704**, 016 (2007); A. Daleo, A. Gehrmann-De Ridder, T. Gehrmann and G. Luisoni, JHEP **1001**, 118 (2010); R. Boughezal, A. Gehrmann-De Ridder and M. Ritzmann, JHEP **1102**, 098 (2011); T. Gehrmann and P. F. Monni, JHEP **1112**, 049 (2011); A. Gehrmann-De Ridder, T. Gehrmann and M. Ritzmann, JHEP **1210**, 047 (2012).
- [27] S. Catani and M. Grazzini, Phys. Rev. Lett. **98** (2007) 222002.
- [28] R. Boughezal, C. Focke, X. Liu and F. Petriello, Phys. Rev. Lett. **115** (2015) 062002; R. Boughezal, X. Liu and F. Petriello, Phys. Rev. D **91** (2015) 094035; J. Gaunt, M. Stahlhofen, F. J. Tackmann and J. R. Walsh, JHEP **1509** (2015) 058.
- [29] M. Czakon, Phys. Lett. B **693** (2010) 259; R. Boughezal, K. Melnikov and F. Petriello, Phys. Rev. D **85** (2012) 034025.
- [30] T. Binoth and G. Heinrich, Nucl. Phys. B **693** (2004) 134; C. Anastasiou, K. Melnikov and F. Petriello, Phys. Rev. D **69** (2004) 076010.
- [31] M. Cacciari, F. A. Dreyer, A. Karlberg, G. P. Salam and G. Zanderighi, Phys. Rev. Lett. **115**, 082002 (2015)
- [32] A. Gehrmann-De Ridder, T. Gehrmann, E. W. N. Glover, A. Huss and T. A. Morgan, Phys. Rev. Lett. **117**, 022001 (2016).
- [33] A. Gehrmann-De Ridder, T. Gehrmann, E. W. N. Glover, A. Huss and T. A. Morgan, arXiv:1601.04569 [hep-ph].
- [34] J. Currie, A. Gehrmann-De Ridder, E. W. N. Glover and J. Pires, JHEP **1401**, 110 (2014).
- [35] M. Cacciari, G. P. Salam and G. Soyez, JHEP **0804**, 063 (2008).
- [36] E. W. N. Glover and J. Pires, JHEP **1006**, 096 (2010); A. Gehrmann-De Ridder, E. W. N. Glover and J. Pires, JHEP **1202**, 141 (2012); A. Gehrmann-De Ridder, T. Gehrmann, E. W. N. Glover and J. Pires, JHEP **1302**, 026 (2013).
- [37] S. Dittmaier, A. Huss and C. Speckner, JHEP **1211**, 095 (2012).
- [38] For Born kinematics (appropriate to LO, virtual NLO and double virtual NNLO contributions) and the hardest jet in any real NLO or double real NNLO event $p_T = p_{T_1}$ exactly. For the subleading jets in real emission contributions this only approximately holds, although the cross section for events with a large hierarchy in p_T between leading and subleading jets is a small contribution to the total. We therefore expect, and observe, the difference between using p_{T_1} and p_T as the RGE evolution variable to be small.